## 58. The Photoelectron Spectrum of [18]Annulene

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## Summary

The PE. spectrum of [18]annulene has been measured and correlated with MOcalculations. The experimental ionization energies can only be explained by computing the electronic states of the cation, that is by taking into account the electron correlation and reorganization in the ionic states. The results allow a discussion of the structure of the neutral molecule; they are consistent with a  $D_{6h}$ point group of symmetry.

Introduction. – In a recent study on the low-temperature UV./VIS. absorption spectrum of [18]annulene we have reviewed the controversy on the structure of this molecule [1]. A structure of  $D_{6h}$  symmetry was proposed on the basis of an X-ray analysis [2], of the UV./VIS. spectrum [3], of NMR. [4], and of thermochemical [5] measurements. On the other hand a structure of  $D_{3h}$  symmetry was predicted by different authors on the basis of various calculations [6–12]. The importance of electron correlation in computations of molecular energies and geometries of large annulenes was studied by one of the authors [13] who concluded that cyclic [4n + 2]- $\pi$ -systems with a twofold axis or a plane through a  $\pi$ -center are stabilized by electron correlation. This finding is now supported by recent MNDOC-computations [14] [15].

In this paper we discuss the PE. spectrum of [18]annulene recorded at 150°. CNDO/S-CI-Calculations [16] [17] are performed for the purpose of interpreting the experimental spectrum. The question of the symmetry of the molecule and its cation is rediscussed.

**Experimental results.** – *PE. Spectrum.* The data were obtained using a *PS-18*-spectrometer from *Perkin-Elmer* (Beaconsfield, England). The calibration was made by introducing small amounts of a mixture of Xe  $({}^{2}P_{1/2}, {}^{2}P_{3/2}$  doublet at 12.130 and 13.436 eV) and Ar  $({}^{2}P_{1/2}, {}^{2}P_{3/2}$  doublet at 15.759 and 15.937 eV) in the target chamber. To avoid oxidation of the [18]annulene the samples were sealed under vacuum in small pyrex ampullas which were opened immediately prior to their introduction into the spectrometer. To get a sufficiently high vapour pressure inside the target chamber the sample holder had

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Fig. 1. Overall PE. spectrum of [18] annulene

to be heated around 150-155°. This temperature corresponds more precisely to the temperature of the metallic insert holding the pyrex ampulla. The actual sample temperature may be substantially lower than 150°. Since [18]annulene was known to thermally decompose in solution above 100° [5] many experiments were carried out to insure that the spectrum reported in *Figure 1* does not contain features from other compounds. One of the decomposition products of [18]annulene is benzene which has a very characteristic PE. band with vibrational structure at 9.25 eV. This band appears about  $1\frac{1}{2}$  h after the introduction of the sample introduction. The benzene band grows slowly with time and at least three new bands, centered around 7.4, 8.3 and 9.5 eV appear. These bands arise from other decomposition intermediates or products like the tetracyclic compound or the benzene [5]. None of these features (and especially the benzene band) are present in the spectrum recorded immediately after the introduction of the sample (*Fig.1*); we thus conclude that it corresponds to the actual PE. spectrum of [18]annulene. The experimentally determined ionization potentials (IP). are reported in *Table 1*.

**Theoretical results.** – Geometry-optimization calculations. According to its crystal structure [2], [18]annulene is nearly planar, has delocalized  $\pi$ -bonds but is

	Band	IP's				
	1	7.23 <sup>b</sup> )				
	2	8.82, 8.99				
	3	10.02				
	4	11.0, 11.5, 12.5, 13.1, 13.5°)				
	5	14.42				
	6	15.34				
	7	16.8				

Table 1. Vertical ionization potentials (IP's, eV)<sup>a</sup>) of [18]annulene

<sup>a</sup>) Maximum of *Frank-Condon* envelope; average of three determinations, ±0.05 or ±0.1 eV.
 <sup>b</sup>) Adiabatic IP.: 6.60 eV; the asymmetry on the high energy side may arise from a *Jahn-Teller* splitting (~0.25 eV).<sup>c</sup>) Very broad band with diffuse maxima.



Fig. 2. Thermal decomposition of [18] annulene as observed by PE. spectroscopy

slightly distorted from an ideal  $D_{6h}$  symmetry: it belongs to the point group  $C_i$ . Neglecting the out-of-plane deviations, a structure belonging to  $C_{2h}$  symmetry (*Fig. 2* in [1]) can be obtained. Moreover, only small corrections are necessary to idealize this latter structure so that it belongs to the  $D_{6h}$  point group of symmetry (*Fig. 3* in [1]). We have performed geometry-optimization calculations using a modified program (OPTMO90)<sup>2</sup>) from *Komornicki & McIver* [18]; the results of this MINDO/2-optimized structure, belonging to the  $D_{3h}$  point group of symmetry, are summarized in *Figure 4* in [1]. It is very similar to the structures found by Allinger et al. [7] and Dewar et al. [8]. For the cation ground state we propose the OPTMO90-[18]optimized dication structure, whereby the  $D_{3h}$ -like  $C_{2v}$  geometry was found to be more stable than the  $D_{6h}$ -like  $C_{2h}$  structure.

CNDO/S-CI-Calculations [16]. The theoretical interpretation of the PE. spectrum was attempted, in a first step, by use of Koopmans' theorem. The ionization

<sup>&</sup>lt;sup>2</sup>) OPTMO90, being suitable to treat large molecules, is a strongly modified version of the original program.

satisfying; especially the second IP. at 8.82, (8.99) eV is not reproduced by this simple approach. We therefore computed the electronic states of the cation by using the  $D_{3h}$ -like  $C_{2v}$ - and  $D_{6h}$ -like  $C_{2h}$ -structures and the program ONDUV99 [19] (Table 2). The second and third IP.'s appear to be due to the  $A_2$ ,  $B_1$  or  $B_g$ ,  $A_u$  states of the  $C_{2\nu}$  and  $C_{2h}$  structures, respectively, which cannot be simply related to one single molecular orbital. In Table 2 the percentages of the one-hole configurations in the computed electronic states of the cation are also given. The first PE, band corresponds to one nearly pure electronic configuration, whereas the second and third bands are due to strongly mixed states. The fourth very broad band with diffuse maxima at 11.0, 11.5, 12.1 and 13.1 eV arises from at least three rather pure (B<sub>2</sub>, B<sub>2</sub>,  $A_1$  or  $B_u$ ,  $B_u$ ,  $A_g$ ) states and from many mixed states. From the data listed in *Table 2*, it is evident that the experimental differences of ionization energies are better reproduced if a  $D_{3h}$ -like structure is assigned to the cation. Our computational

Nr	CNDO/S-Orbitals						CNDO/S-CI-States					Exp.		
	Energy	$D_{3h}, C_{2v}$ T	ype	Energy	'D <sub>6h</sub> '	C <sub>2h</sub> T	уре	Energy	$C_{2v}$	Dne- nole %	Energy	$C_{2h}$	Dne- hole %	
1	7.91 8.00	e" b <sub>1</sub> a <sub>2</sub>	π π	7.26 7.30	e <sub>2u</sub>	a <sub>u</sub> a <sub>u</sub>	π π	7.85 8.44	B <sub>1</sub> A <sub>2</sub>	98 87	7.74 7.94	A <sub>u</sub> A <sub>u</sub>	96 94	7.23
2								9.75 9.91 9.96	$\begin{array}{c} A_2 \\ B_1 \\ A_2 \end{array}$	2 38 44	9.19 9.31 9.46	Bg Bg Bg	2 6 31	8.82 8.99
3	10.16 10.21	$a_1^{"} a_2$ $a_2^{"} b_1$	π	9.91 10.11	b <sub>2g</sub> b <sub>1g</sub>	b <sub>g</sub> b <sub>g</sub>	π	10.70 11.05 11.18	B <sub>1</sub> A <sub>2</sub> B <sub>1</sub>	11 50 46	9.61 10.89 10.96 11.29 11.44 11.53	$ \begin{array}{c}  B_g \\  A_u \\  A_u \\  B_g \\  B_g \\  B_g \end{array} $	31 20 18 17 57 39	10.02
4	11.30	aź b <sub>2</sub>	σ	11.98	b <sub>2u</sub>	b <sub>u</sub>	σ	11.49 11.55 11.60 11.64	$\begin{array}{c} A_2 \\ B_1 \\ B_2 \\ A_2 \end{array}$	12 22 99 18	11.74 11.78 11.81 12.50	$ \begin{array}{c}         B_{g} \\         A_{u} \\         B_{g} \\         A_{u} \end{array} $	2 3 1 0	11.00 11.50
	11.74 11.80	e' a <sub>1</sub> b <sub>2</sub>	σ σ	12.15 12.26	e <sub>2u</sub>	a <sub>u</sub> a <sub>u</sub>	π π	11.97 12.21 12.47 12.56	$     B_1 \\     B_2 \\     A_1 \\     A_2 $	0 88 95 8	12.52 12.67 12.83 12.85	$A_u \\ B_g \\ B_u \\ A_u$	0 3 100 0	12.50
	12.27 12.29	e" a <sub>2</sub> b <sub>1</sub>	π π	12.27 12.32	e <sub>2g</sub>	a <sub>g</sub> b <sub>u</sub>	σ σ	12.63 13.06 13.11 13.14 13.17 13.22	$B_1$ $A_2$ $B_1$ $B_1$ $A_2$ $A_2$	17 47 51 2 15 2	12.85 12.95 13.12 13.13 13.15 13.19	$B_g$ $A_u$ $B_u$ $A_u$ $A_u$ $A_g$	2 99 44 39 99	13.10
								13.42 13.43 13.56	A <sub>2</sub> B <sub>1</sub> A <sub>1</sub>	0 1 0				

Table 2. CNDO/S-orbital energies of [18]annulene and CNDO/S-CI-computed states of its cation (eV)

D <sub>6h</sub>	MNDO	MNDOC	D <sub>3h</sub>	MNDO	MNDOC	
E <sub>211</sub>	6.99	7.15	E″	7.94	8.04	
B <sub>2</sub> ,	8.64	9.07	$A_2''$	9.29	9.62	
Bie	8.73	9.14	$A_1^{\tilde{\prime}\prime}$	9.36	9.69	
$E_{2u}$	10.24	10.78	E″	10.61	11.15	

Table 3. MNDO- [20] and MNDOC- [15] computed states of the cation of [18] annulene

results in *Table 3* may be compared with those of the MNDO- [20] and MNDOC- [15] calculations of *Thiel* [21].

**Conclusions.** – On the basis of many experimental facts and their interpretation [2–5] we are confident that [18]annulene has a structure with  $D_{6h}$ -symmetry. The computational preference for a  $D_{3h}$ -symmetry is almost certainly due to the neglect of the correlation energy. On the other hand it is rather sure that the cation of this molecule (no (4n+2)-system, where the correlation is important) has localized  $\pi$ -bonds and is best described by a distorted  $D_{3h}$ -structure, often proposed for the neutral compound.

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